

Note From last day...

$$\gamma^2 = -\omega^2 \mu \epsilon_c = -\omega^2 \mu \left(\epsilon - j\frac{\sigma}{\omega} \right)$$

$$\nabla^2 \vec{E}_s - \gamma^2 \vec{E}_s = 0$$

Plane waves: defined in a direction \hat{n}

$$\vec{E}_s = \vec{E}_0 e^{\pm \gamma \hat{n} \cdot \vec{r}}$$

$$\gamma = \alpha + j\beta$$

↳ possibly complex constant.

Special case: $\hat{n} = \hat{z}$ $\hat{n} \cdot \vec{r} = z$

\vec{E}_0 is real

$$\vec{E}(z, t) = \text{Re} \left[\vec{E}_0 e^{-\alpha z} e^{-j\beta z} e^{j\omega t} \right]$$

$$= \vec{E}_0 e^{-\alpha z} \cos(\omega t - \beta z)$$

$$\vec{H}(z, t) = \vec{H}_0 e^{-\alpha z} \cos(\omega t - \beta z)$$

Show that $E_{sz} = 0$

$$\nabla \times \vec{H}_s = j\omega \epsilon_c \vec{E}_s$$

↳ z component of this

$$\frac{\partial}{\partial x} H_{sy} - \frac{\partial}{\partial y} H_{sx} = 0$$

$$\Rightarrow j\omega \epsilon_c E_{sz} = 0$$

$$\therefore E_{sz} = 0$$

Thus if $\hat{n} = \hat{z}$, \vec{E} is \perp to \hat{n}

Now to show that the magnetic field is also \perp to the \hat{n}

So: $\vec{E}_0 = \hat{x} E_{0x} + \hat{y} E_{0y}$

$$\vec{\nabla} \times \vec{E}_s = -j\omega\mu \vec{H}_s$$

$$= \begin{bmatrix} \hat{x} & \hat{y} & \hat{z} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ E_{0x}e^{-\delta z} & E_{0y}e^{-\delta z} & 0 \end{bmatrix}$$

$$= -\delta(-\hat{x} E_{0y} + \hat{y} E_{0x})e^{-\delta z}$$

\perp to $\hat{n} = \hat{z}$

$$= -j\omega\mu \vec{H}_s$$

$\therefore \vec{H}_s$ is \perp to \hat{n}

$$\vec{H}_s = (\hat{x} H_{0x} + \hat{y} H_{0y})e^{-\delta z} \quad \left. \begin{array}{l} \text{by} \\ \text{comparing} \end{array} \right\}$$

$$\therefore H_{0x} = \frac{-j}{\omega\mu} (-\delta) E_{0y}$$

$$H_{0y} = \frac{j}{\omega\mu} (-\delta) E_{0x}$$

$$\vec{H}_s \cdot \vec{E}_s = (\vec{E}_0 \cdot \vec{H}_0) e^{-2\delta z}$$

$$= 0$$

Summary

$$\vec{E} \cdot \hat{n} = 0$$

$$\vec{H} \cdot \hat{n} = 0$$

$$\vec{H} \cdot \vec{E} = 0$$

plane waves

Solve for $\gamma = \alpha + j\beta$ in terms of μ, ϵ, σ .

$$\begin{aligned}\gamma^2 &= -\omega^2 \mu \epsilon \left(1 - \frac{j\sigma}{\omega}\right) \\ &= (-1) \omega^2 \mu \epsilon \left(1 - \frac{j\sigma}{\omega}\right)\end{aligned}$$

$$\gamma = j\omega \sqrt{\mu \epsilon} \left(1 - \frac{j\sigma}{\omega}\right)^{1/2}$$

$$\gamma = \alpha + j\beta \quad ; \quad \gamma^2 = \alpha^2 - \beta^2 + 2j\alpha\beta$$

$$\alpha^2 = \beta^2 = -\omega^2 \mu \epsilon$$

$$2j\alpha\beta = -j\omega \mu \epsilon \sigma$$

$$\therefore \beta = \frac{\omega \epsilon \sigma \mu}{2\alpha}$$

$$\alpha^2 - \beta^2 = \alpha^2 - \left(\frac{\omega^2 \epsilon^2 \sigma^2 \mu^2}{2\alpha^2}\right)$$

Oh... he fucked this right up... check the text... plus these are given on cheat sheet.

$$\alpha = \sqrt{\frac{\omega^2 \mu \epsilon}{2}} \left[\left(1 + \frac{\sigma^2}{\omega^2 \epsilon^2}\right)^{1/2} - 1 \right]$$

$$\beta = \sqrt{\frac{\omega^2 \mu \epsilon}{2}} \left[\left(1 + \frac{\sigma^2}{\omega^2 \epsilon^2}\right)^{1/2} + 1 \right]$$

Note: $\sigma = 0$ then $\alpha = 0$
 $\beta = \sqrt{\omega^2 \mu \epsilon}$

Assume: $\hat{n} = \hat{z}$

$$\vec{E}_0 = \hat{x} |E_{0x}| e^{j\theta_x}$$

$$\vec{H}_0 = \hat{y} |H_{0y}| e^{j\theta_y}$$

$$\vec{E}(z,t) = \text{Re} \left[\hat{x} |E_{0x}| e^{j\theta_x} e^{-\alpha z} e^{-j\beta z} e^{j\omega t} \right]$$

$$\{ = \hat{x} |E_{0x}| e^{-\alpha z} \cos(\omega t - \beta z + \theta_x)$$

right traveling

$\delta = \frac{1}{\alpha}$: characteristic of how deep the wave propagates.

$$\{ \vec{E}(z,t) = \hat{x} |E_{0x}| e^{\alpha z} \cos(\omega t + \beta z + \theta_x)$$

left traveling.

$$u = \pm \omega / \beta$$

wavelength: distance between two peaks assuming $\alpha = 0$.

$$\lambda = \frac{2\pi}{\beta}$$

Units of \vec{E} : V/m

Units of \vec{H} : Amp/m

$$\frac{|\vec{E}|}{|\vec{H}|} = \frac{V}{I}$$

note there is nothing to guarantee that both vectors will reach their peak at the same time.

Power density: $|\vec{E}| |\vec{H}|$